# Numerical Uniqueness and Existence Theorem for Solution of Lippmann-Schwinger Equation 

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## Abstract-

In this paper, numerical uniqueness and existence theorem is presented for solution of LippmannSchwinger equation for sound scattering problem. A remarkable feature of this theorem is that the sufficient condition shown in this theorem can be evaluated by verified numerical computations.

## 1. Introduction

In this paper, we are concerned with the following scattering problem for the Helmholtz equation in the inhomogeneous media: find $u: \mathbb{R}^{2} \rightarrow \mathbb{C}$ such that

$$
\begin{align*}
& \triangle u(x)+\kappa^{2} b(x) u(x)=0, x \in \mathbb{R}^{2}  \tag{1}\\
& u=u^{i}+u^{s},  \tag{2}\\
& \lim _{r=|x| \rightarrow \infty} r^{1 / 2}\left(\frac{\partial u^{s}}{\partial r}-\mathbf{i} \kappa u^{s}\right)=0 . \tag{3}
\end{align*}
$$

Here, $\mathbf{i}=\sqrt{-1}$ and $\kappa$ is the wave number of the incidence wave. It is well known that this problem is given as a mathematical modeling of a two dimensional sound scattering problem [1].

We here assume the following:
Assumption 1 The refractive index $b: \mathbb{R}^{2} \rightarrow \mathbb{C}$ is $a$ given smooth function and

$$
\begin{equation*}
a(x)=b(x)-1 \tag{4}
\end{equation*}
$$

has a compact support such that

$$
\begin{equation*}
\operatorname{supp} a \subset \bar{B}_{\rho}, \quad B_{\rho}=\left\{x \in \mathbb{R}^{2}:|x|<\rho\right\} . \tag{5}
\end{equation*}
$$

Furthermore, we assume that

$$
\begin{equation*}
a \in W^{\mu, 2}\left(\mathbb{R}^{2}\right) \tag{6}
\end{equation*}
$$

with $\mu>1$.
Here, the Sobolev space $W^{\mu, 2}\left(\mathbb{R}^{2}\right)$ consists of functions $w \in L^{2}\left(\mathbb{R}^{2}\right)$ such that

$$
\begin{equation*}
\int_{\mathbb{R}^{2}}\left(1+|\xi|^{2}\right)^{\mu}|\hat{w}(\xi)|^{2} d \xi<\infty \tag{7}
\end{equation*}
$$

where

$$
\begin{equation*}
\hat{w}(\xi)=\int_{\mathbb{R}^{2}} \mathbf{e}^{-\mathbf{i} x \cdot \xi} w(x) d x, \xi \in \mathbb{R}^{2} \tag{8}
\end{equation*}
$$

is the Fourier transform of $w$.
The incidence wave is assumed to be

$$
\begin{equation*}
u^{i}(x)=\exp (\mathbf{i} \kappa d \cdot x) \tag{9}
\end{equation*}
$$

with a fixed $d \in \mathbb{R}^{2},|d|=1$. The incidence wave $u^{i}$ is a plane wave solution of the Helmholtz equation

$$
\begin{equation*}
\triangle u(x)+\kappa^{2} u(x)=0 \tag{10}
\end{equation*}
$$

The scattering wave $u^{s}$ for this incidence wave is assumed to satisfy the Sommerfeld radiation condition (3).

It is known [1] that the problem defined by (1)-(3) is equivalent to the Lippmann-Schwinger equation

$$
\begin{equation*}
u(x)=\kappa^{2} \int_{\mathbb{R}^{2}} E_{\kappa}(x-y) a(y) u(y) d y+u^{i}(x), x \in \mathbb{R}^{2} \tag{11}
\end{equation*}
$$

where $E_{\kappa}(x)=(\mathbf{i} / 4) H_{0}^{(1)}(\kappa|x|)$. Here, $H_{0}^{(1)}$ is the Hankel function of the first kind of order zero. A fast solver for this equation has been proposed and studied mathematically in detail by Vainikko [2]. The monograph Saranen-Vainikko [3] has described Vainikko's theory in detail with its mathematical background. According to [2] and [3], first, we somewhat simplify and generalize (11). If we scale the independent variables $x$ and $y$ by $\tilde{x}=\kappa x$ and $\tilde{y}=\kappa y$, respectively, without loss of generality, we can assume that $\kappa=1$. Further, instead of $u^{i}$, which is a solution of homogenized Helmholtz equation with $b(x)=1$ for all $x \in \mathbb{R}^{2}$, we considier a general function $f$. We assume the following:

Assumption $2 f \in W_{\mathrm{loc}}^{\mu, 2}\left(\mathbb{R}^{2}\right), \mu>1$.
Let us recall that $a \in W^{\mu, 2}\left(\mathbb{R}^{2}\right)$ with supp $a \subset \bar{B}_{\rho}=$ $\left\{x \in \mathbb{R}^{2}:|x| \leqq \rho\right\}$. Then, the problem can be formulated as: find $u \in C\left(\bar{B}_{\rho}\right)$ satisfying the following equation

$$
\begin{equation*}
u(x)=\int_{B_{\rho}} E_{1}(x-y) a(y) u(y) d y+f(x), x \in B_{\rho} . \tag{12}
\end{equation*}
$$

By the Sobolev embedding theorem, $W_{\text {loc }}^{\mu, 2} \subset C\left(\mathbb{R}^{2}\right)$ for $\mu>1$. Multiplying both sides of (12) by $a(x)$, we can rewrite the equation with respect to $v(x)=a(x) u(x)$ as an unknown function:

$$
\begin{equation*}
v(x)=a(x) \int_{B_{\rho}} E_{1}(x-y) v(y) d y+a(x) f(x), x \in B_{\rho} . \tag{13}
\end{equation*}
$$

A crucial observation is that, for $x \in \bar{B}_{\rho}$, only the values from $\bar{B}_{2 \rho}$ of $E_{1}(x)$ are involved in the integral, i.e., changing $E_{1}$ outside this ball, the solution $v(x)$ does not change in $\bar{B}_{\rho}$. We exploit this observation and redefine the kernel $E_{1}$ in $G_{R} \backslash \bar{B}_{R}$ where

$$
G_{R}=\left\{x=\left(x_{1}, x_{2}\right) \in \mathbb{R}^{2}:\left|x_{1}\right|<R,\left|x_{2}\right|<R\right\}
$$

is an open box and $R>2 \rho$ is a parameter. Namely, we define

$$
\mathcal{K}(x)= \begin{cases}E_{1}(x)=\frac{\mathbf{i}}{4} H_{0}^{(1)}(|x|), & |x| \leqq R  \tag{14}\\ 0, & x \in G_{R} \backslash \bar{B}_{R}\end{cases}
$$

After that we extend $\mathcal{K}, a, a f$ and $v$ from $G_{R}$ to $\mathbb{R}^{2}$ as $2 R$-periodic functions with respect to $x_{1}$ and $x_{2}$. Thus, we have arrived at the biperiodic integral equation

$$
\begin{equation*}
v(x)=a(x) \int_{G_{R}} \mathcal{K}(x-y) v(y) d y+a(x) f(x) \tag{15}
\end{equation*}
$$

which is equivalent to (12).
The main theorem in [2] and [3] is the following:
Theorem 1 (Vainikko) Assume that functions a and $f$ satisfies Assumptions 1, 2 mentioned above and the homogeneous equation corresponding to (15) has in $H^{0}$ only the trivial solution. Then (15) has a unique solution $v \in H^{\mu}$ has a unique solution $v_{N} \in \mathcal{T}_{N}$ for sufficiently large $N$, and

$$
\begin{equation*}
\left\|v_{N}-v\right\|_{\lambda} \leqq c N^{\lambda-\mu}\|v\|_{\mu}, \quad(0 \leqq \lambda \leqq \mu) \tag{16}
\end{equation*}
$$

In [1], the existence of a unique solution in a function space of continuous functions is proved for the Lippmann-Schwinger equation in $\mathbb{R}^{3}$ using the unique continuation principle.

In this paper, for the sake of simplicity, we treat the case of $\mu \geqq 3 / 2$ and present Theorem giving a sufficient condition of guaranteeing that the homogeneous equation corresponding to (15) has in $H^{0}$ only the trivial solution. Thus, it can be seen as an another uniqueness theorem. A remarkable feature of this theorem is that the sufficient condition shown in this theorem can be evaluated by verified numerical computations. Furthermore, if such a sufficient condition holds, it also gives an upper bound of $\left\|(I-a K)^{-1}\right\|_{\mathcal{L}\left(H^{0}\right)}$ and a tight error bound between the exact solution $v$ and an approximate solution $\tilde{v}$ which is generated by computer. As a result, this paper is to present a numerical uniqueness and existence theorem for (15), which asserts the existence of a unique solution around an approximate solution computed by numerical calculation.

## 2. Verification Theory

To present a numerical existence theorem of (15), we first prove several lemma.
Lemma 1 Let $a \in H^{\mu}$ and $\mu \geqq \frac{3}{2}$. Then,

$$
\begin{equation*}
\|a K\|_{\mathcal{L}\left(H^{0}, H^{\frac{3}{2}}\right)} \leqq c_{\frac{3}{2}} c_{R}\|a\|_{\frac{3}{2}} . \tag{17}
\end{equation*}
$$

Proof It is know that the following holds:

$$
\begin{equation*}
\|a K v\|_{\frac{3}{2}} \leqq c_{\frac{3}{2}}\|a\|_{\frac{3}{2}}\|K v\|_{\frac{3}{2}} \tag{18}
\end{equation*}
$$

Further, we have

$$
\begin{align*}
\|K v\|_{\frac{3}{2}} & =\left\|\int_{G_{R}} \mathcal{K}(x-y) v(y) d y\right\|_{\frac{3}{2}} \\
& =\sqrt{\sum_{j \in \mathbb{Z}^{2}} \underline{j}^{3}|\hat{F}(j)|^{2}} \\
& =\sqrt{\sum_{j \in \mathbb{Z}^{2}} \underline{j}^{3}|\hat{K}(j)|^{2}|\hat{v}(j)|^{2}} \\
& \leqq \sqrt{\sum_{j \in \mathbb{Z}^{2}} c_{R}^{2}|\hat{v}(j)|^{2}} \\
& =c_{R}\|v\|_{0} . \tag{19}
\end{align*}
$$

Here, we have put

$$
\begin{equation*}
F(x)=\int_{G_{R}} \mathcal{K}(x-y) v(y) d y \tag{20}
\end{equation*}
$$

Further, we have used

$$
\begin{equation*}
|\mathcal{K}(j)| \leqq c_{R} \underline{j}^{-\frac{3}{2}} \tag{21}
\end{equation*}
$$

QED
Lemma 2 Let $a \in H^{\mu}$ and $\mu \geqq \frac{3}{2}$. Then,

$$
\begin{equation*}
\left\|Q_{N}(a K)\right\|_{\mathcal{L}\left(H^{0}, \mathcal{T}_{N}\right)} \leqq c_{\frac{3}{2}} c_{R}\|a\|_{\frac{3}{2}} . \tag{22}
\end{equation*}
$$

Proof

$$
\begin{align*}
\left\|Q_{N}(a K v)\right\|_{0} & \leqq\|a K v\|_{\frac{3}{2}} \\
& =c_{\frac{3}{2}} c_{R}\|a\|_{\frac{3}{2}}\|v\|_{0} \tag{23}
\end{align*}
$$

QED
Theorem 2 Let $a \in H^{\mu}$ and $\mu \geqq \frac{3}{2}$. Let $k=$ $c_{\frac{3}{2}} c_{R}\|a\|_{\frac{3}{2}}$ and $c_{N}=c_{0, \frac{3}{2}} N^{-\frac{3}{2}}$. If

$$
\begin{equation*}
\left\|\left(I-Q_{N} a K\right)^{-1}\right\|_{\mathcal{L}\left(\mathcal{T}_{N}\right)} \leqq M \tag{24}
\end{equation*}
$$

and

$$
\begin{equation*}
c_{N} k(1+M k)<1 \tag{25}
\end{equation*}
$$

the operator $I-a K$ is invertible and the following is satisfied:

$$
\begin{equation*}
\left\|(I-a K)^{-1}\right\|_{\mathcal{L}\left(H^{0}\right)} \leqq C_{M} \tag{26}
\end{equation*}
$$

Here,

$$
\begin{equation*}
C_{M}=\frac{1+M k}{1-c_{N} k(1+M k)} \tag{27}
\end{equation*}
$$

From Therorem 2, it follows the following theorem:
Theorem 3 Let $a \in H^{\mu}$ and $\mu \geqq \frac{3}{2}$. Let $\tilde{v} \in \mathcal{T}_{N}$ be an approximate solution of

$$
\begin{equation*}
(I-a K) v=a f \tag{28}
\end{equation*}
$$

If

$$
\begin{equation*}
\left\|\left(I-Q_{N} a K\right)^{-1}\right\|_{\mathcal{L}\left(\mathcal{T}_{N}\right)} \leqq M \tag{29}
\end{equation*}
$$

and

$$
\begin{equation*}
c_{N} k(1+M k)<1 \tag{30}
\end{equation*}
$$

hold, then (28) has a unique solution $v^{*}$ satisfying

$$
\begin{equation*}
\left\|\tilde{v}-v^{*}\right\|_{H^{0}} \leqq C_{M}\|(I-a K) \tilde{v}-a f\|_{H^{0}} . \tag{31}
\end{equation*}
$$

Furthermore, $v^{*} \in H^{\frac{3}{2}}$.

## References

[1] David Colton and Rainer Kress: Inverse Acoustic and Electromagnetic Scattering Theory (Second Edition), Applied Mathematical Sciences 93, Springer (1998).
[2] Gennadi Vainikko:" Fast Solvers of the LippmannSchwinger Equation", Helsinki University of Technology, Institute of Mathematics Research Reports, A387 (1997).
[3] Jukka Saranen and Genneadi Vainikko: Periodic Integral and Pseudodifferential Equations with Numerical Approximation, Springer Monographs in Mathematics, Springer (2002).

