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# A Method for the Computation of Border Collision Bifurcation Point in a Piecewise Linear System with Interrupted Characteristics 

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#### Abstract

This paper proposes a simple method for analyzing the border-collision bifurcation point in a piecewise linear interrupted system. First, we show the system dynamics, and then we derive the Poincaré map. Next, we explain the method for calculating the border-collision bifurcation point based on the Poincaré map approach. Finally, we apply the method to the current mode controlled buck-boost converter. The analysis result will show the validity of the method.


## 1. Introduction

The piecewise smooth systems (PWS) have a interrupted characteristic due to the complex switching actions. There are many papers that have studied the bifurcation phenomena in the PWS. We know that it is important to derive the bifurcation sets for understanding the qualitative property of the PWS. The calculation method for the local bifurcations, such as the period-doubling bifurcation and Neimark-Sacker bifurcation, have been studied. On the other hand, a new non-local bifurcation phenomenon, called the border-collision bifurcation, has been discovered in the PWS [1]. Many papers have studied mechanism of the border-collision bifurcation [2,3]. In addition, some papers have proposed the calculation method of the bordercollision bifurcation point based on the Poincaré map approach [4].

The sampled data model of the interrupted electric circuit, whose switch is dependent on the state and a periodic interval, often constructs the PWS. It means that the bordercollision bifurcation can occur in the interrupted electric circuit, and indeed Refs. [5,6] have reported appearance of the border-collision bifurcation. Of course there is an algorithm for calculating the bifurcation point of the bordercollision [7]. In theory, we can calculate the bifurcation point by using above algorithm even if the system has nonlinear characteristics and is high-dimensional. But, the algorithm [7] is complicated because it needs to analytically derive the composite Poincaré map and its derivative. For
this reason, it is valuable to construct the algorithm with no above complicated computation process.

In this paper, we aim to construct a simple algorithm for calculating the bifurcation point of the border-collision with the help of the mathematical tool MAPLE; the MAPLE will help to derive the composite Poincaré map and its derivative even if the system is high-dimensional [8]. First, we show the algorithm based on the Poincaré map approach. Next, we apply the method to the current mode controlled buck-boost converter. It will be possible to calculate the bifurcation point of the high-dimensional system by using the proposed algorithm. But, we apply the algorithm for the two-dimensional buck-boost converter as the first step to confirm the validity of the algorithm. Finally, we show that the algorithm can calculate the bifurcation point of the border-collision.

## 2. Bifurcation analysis

### 2.1. System description

We consider the following differential equations:

$$
\begin{align*}
& \left\{\begin{array}{l}
\frac{d x}{d t}=f_{\mathrm{a}}\left(x, y, \lambda_{\mathrm{a}}, \lambda_{\mathrm{c}}\right) \\
\frac{d y}{d t}=g_{\mathrm{a}}\left(x, y, \lambda_{\mathrm{a}}, \lambda_{\mathrm{c}}\right)
\end{array},\right. \text { System-a, }  \tag{1}\\
& \left\{\begin{array}{l}
\frac{d x}{d t}=f_{\mathrm{b}}\left(x, y, \lambda_{\mathrm{b}}, \lambda_{\mathrm{c}}\right) \\
\frac{d y}{d t}=g_{\mathrm{b}}\left(x, y, \lambda_{\mathrm{b}}, \lambda_{\mathrm{c}}\right)
\end{array},\right. \text { System-b, } \tag{2}
\end{align*}
$$

where $\lambda_{\mathrm{a}}, \lambda_{\mathrm{b}} \in \boldsymbol{R}^{r}$ and $\lambda_{\mathrm{c}} \in \boldsymbol{R}^{s}$ are the system parameters. Let the exact solutions of Eqs. (1) and (2) be:

$$
\begin{align*}
& \left\{\begin{array}{l}
x(t)=\varphi_{\mathrm{a}}\left(t ; x, y, \lambda_{\mathrm{a}}, \lambda_{\mathrm{c}}\right) \\
y(t)=\phi_{\mathrm{a}}\left(t ; x, y, \lambda_{\mathrm{a}}, \lambda_{\mathrm{c}}\right)
\end{array},\right. \text { System-a, }  \tag{3}\\
& \left\{\begin{array}{l}
x(t)=\varphi_{\mathrm{b}}\left(t ; x, y, \lambda_{\mathrm{b}}, \lambda_{\mathrm{c}}\right) \\
y(t)=\phi_{\mathrm{b}}\left(t ; x, y, \lambda_{\mathrm{b}}, \lambda_{\mathrm{c}}\right)
\end{array}\right. \text {, System-b. } \tag{4}
\end{align*}
$$

A local section $\Pi$ in the state space is described as follows:

$$
\begin{equation*}
\Pi=\left\{\boldsymbol{x} \in \boldsymbol{R}^{2} \mid q(y)=0, q: \boldsymbol{R}^{2} \rightarrow \boldsymbol{R}\right\} . \tag{5}
\end{equation*}
$$

Note that the local section $\Pi$ is a scalar function whose is dependent on the state variable $y$.

Figure 1 shows a conceptual diagram for understanding behavior of the waveforms. Let the initial waveform at $t=$ $k T$ be $\boldsymbol{x}_{k}=\boldsymbol{x}(k T)=\left(x_{k}, y_{k}\right)^{\top}$ with System-a. We assume that there are two types of the waveform behaviors during a time interval $T$. The one keeps System-a during a time interval $T$. The other one keeps System-a during a time $\tau_{\mathrm{a}}\left(\boldsymbol{x}_{k}\right)$, and then the System keeps System-b during a time $T-\tau_{\mathrm{a}}\left(\boldsymbol{x}_{k}\right)$.

At the next stage, we derive the Poincaré map for the following analysis. A map $M_{\mathrm{A} k}$ is defined by the following equation if $\tau_{\mathrm{a}} \geq T$ is satisfied.

$$
\begin{align*}
& M_{\mathrm{A} k}: \boldsymbol{R}^{2} \rightarrow \boldsymbol{R}^{2} \\
& \qquad\left\{\begin{array}{l}
x_{k} \longmapsto x_{k+1}=\varphi_{\mathrm{a}}\left(T, x_{k}, y_{k}, \lambda_{\mathrm{a}}, \lambda_{\mathrm{c}}\right) \\
y_{k} \longmapsto y_{k+1}=\phi_{\mathrm{a}}\left(T, x_{k}, y_{k}, \lambda_{\mathrm{a}}, \lambda_{\mathrm{c}}\right)
\end{array} .\right. \tag{6}
\end{align*}
$$

Likewise, a map $M_{\mathrm{B} k}$ is defined by the following equation if $\tau_{a}<T$ is satisfied.

$$
\begin{align*}
& M_{\mathrm{B} k}: \boldsymbol{R}^{2} \\
& \boldsymbol{x}_{k} \mapsto \boldsymbol{R}^{2}  \tag{7}\\
& \longmapsto \boldsymbol{x}_{k+1}=M_{\mathrm{b}} \circ M_{\mathrm{a}} .
\end{align*}
$$

The maps $M_{\mathrm{a}}$ and $M_{\mathrm{b}}$ is defined as follows:

$$
\begin{align*}
& M_{\mathrm{a}}: \boldsymbol{R}^{2} \rightarrow \Pi \\
& \qquad\left\{\begin{array}{l}
x_{k} \longmapsto x_{q}=\varphi_{\mathrm{a}}\left(\tau_{\mathrm{a}}, x_{k}, y_{k}, \lambda_{\mathrm{a}}, \lambda_{\mathrm{c}}\right) \\
y_{k} \longmapsto y_{q}=\phi_{\mathrm{a}}\left(\tau_{\mathrm{a}}, x_{k}, y_{k}, \lambda_{\mathrm{a}}, \lambda_{\mathrm{c}}\right)
\end{array},\right. \tag{8}
\end{align*}
$$

$$
\begin{align*}
& M_{\mathrm{b}}: \Pi \rightarrow \boldsymbol{R}^{2} \\
& \qquad\left\{\begin{array}{l}
x_{q} \longmapsto x_{k+1}=\varphi_{\mathrm{b}}\left(T-\tau_{\mathrm{a}}, x_{q}, y_{q}, \boldsymbol{\lambda}_{\mathrm{b}}, \lambda_{\mathrm{c}}\right) \\
y_{q} \longmapsto y_{k+1}=\phi_{\mathrm{b}}\left(T-\tau_{\mathrm{a}}, x_{q}, y_{q}, \lambda_{\mathrm{b}}, \lambda_{\mathrm{c}}\right)
\end{array} .\right. \tag{9}
\end{align*}
$$



Figure 1: Example of behavior.

Based on Eqs. (6) and (7), let the exact Poincaré map be

$$
\begin{align*}
\boldsymbol{x}_{k+1} & =\left[\begin{array}{c}
F\left(1, x_{k}, y_{k}, \lambda_{\mathrm{a}}, \lambda_{\mathrm{b}}, \lambda_{\mathrm{c}}\right) \\
G\left(1, x_{k}, y_{k}, \lambda_{\mathrm{a}}, \lambda_{\mathrm{b}}, \lambda_{\mathrm{c}}\right)
\end{array}\right] \\
& =\left\{\begin{array}{l}
{\left[\begin{array}{c}
F_{\mathrm{A}}\left(x_{k}, y_{k}, \lambda_{\mathrm{a}}, \lambda_{\mathrm{c}}\right) \\
G_{\mathrm{A}}\left(x_{k}, y_{k}, \lambda_{\mathrm{a}}, \lambda_{\mathrm{c}}\right)
\end{array}\right], \quad y_{k} \leq D} \\
{\left[\begin{array}{c}
F_{\mathrm{B}}\left(x_{k}, y_{k}, \lambda_{\mathrm{a}}, \lambda_{\mathrm{b}}, \lambda_{\mathrm{c}}\right) \\
G_{\mathrm{B}}\left(x_{k}, y_{k}, \lambda_{\mathrm{a}}, \lambda_{\mathrm{b}}, \lambda_{\mathrm{c}}\right)
\end{array}\right], y_{k}>D}
\end{array}\right. \tag{10}
\end{align*}
$$

where a border $D$ is

$$
\begin{equation*}
D=\phi_{\mathrm{a}}\left(-T, x_{k+1}, y_{k+1}, \lambda_{\mathrm{a}}, \lambda_{\mathrm{c}}\right) . \tag{11}
\end{equation*}
$$

### 2.2. Algorithm for calculating bifurcation point

The conditional equation for the period- $p$ solution is defined as follows:

$$
\left[\begin{array}{c}
F\left(p, x_{k}, y_{k}, \lambda_{\mathrm{a}}, \lambda_{\mathrm{b}}, \lambda_{\mathrm{c}}\right)-x_{k}  \tag{12}\\
G\left(p, x_{k}, y_{k}, \lambda_{\mathrm{a}}, \lambda_{\mathrm{b}}, \lambda_{\mathrm{c}}\right)-y_{k}
\end{array}\right]=\mathbf{0} .
$$

We assume that the border-collision bifurcation occurs when a solution $y_{k}$ just on the border $D$. This situation is expressed as follows:

$$
\begin{equation*}
y_{k}-D=0 . \tag{13}
\end{equation*}
$$

Based on Eqs. (12) and (13), we get

$$
P\left(x_{k}, y_{k}, \lambda_{\mathrm{c} 1}\right)=\left[\begin{array}{c}
F\left(p, x_{k}, y_{k}, \lambda_{\mathrm{c} 1}\right)-x_{k}  \tag{14}\\
G\left(p, x_{k}, y_{k}, \lambda_{\mathrm{c} 1}\right)-y_{k} \\
y_{k}-D
\end{array}\right]=\mathbf{0},
$$

where a parameter $\lambda_{\mathrm{c} 1}$ belongs to a parameter $\lambda_{\mathrm{c}}$. Although the previous algorithm numerically calculates Eq. (14) ??, the proposed algorithm can derive Eq. (14) with the exact solution. But, we need the numerical computation then for calculating the bifurcation point of border collision. The Jacobian matrix is derived based on Eq. (14). Specifically, we get

$$
D P\left(x_{k}, y_{k}, \lambda_{\mathrm{c} 1}\right)=\left[\begin{array}{ccc}
\frac{\partial F}{\partial x_{k}}-1 & \frac{\partial F}{\partial y_{k}} & \frac{\partial F}{\partial \lambda_{\mathrm{c} 1}}  \tag{15}\\
\frac{\partial G}{\partial x_{k}} & \frac{\partial G}{\partial y_{k}}-1 & \frac{\partial G}{\partial \lambda_{\mathrm{c} 1}} \\
0 & 1 & \frac{\partial \phi_{\mathrm{a}}}{\partial \lambda_{\mathrm{c} 1}}
\end{array}\right] .
$$

## 3. Example of the application

We apply the method to the current mode controlled buck-boost converter (see Fig. 2). Note that we only consider the continuous conduction mode for the sake of the simplicity. Moreover, we use the dimensionless values: $y=\sqrt{L} i, x=\sqrt{C} v, \tau=t / \sqrt{L C}, \sigma=\sqrt{L / C} /(2 R)$, and $B=\sqrt{C} E$. The circuit equations are derived as follows:

$$
\left\{\begin{array}{l}
\frac{d x}{d \tau}=-2 \sigma x  \tag{16}\\
\frac{d y}{d \tau}=B
\end{array}, \mathrm{SW}: \mathrm{ON}\right.
$$

$$
\left\{\begin{array}{l}
\frac{d x}{d \tau}=y-2 \sigma x  \tag{17}\\
\frac{d y}{d \tau}=-x
\end{array}, \mathrm{SW}: \mathrm{OFF}\right.
$$

Figure 3 shows the behavior of the waveforms. The Poincaré map is derived with the exact solution.

$$
\begin{align*}
& {\left[\begin{array}{l}
x_{k+1} \\
y_{k+1}
\end{array}\right]} \\
& =\left\{\begin{array}{c}
{\left[\begin{array}{c}
x_{k} e^{2 \alpha T} \\
y_{k}+B T
\end{array}\right],} \\
{\left[\begin{array}{c}
e^{\alpha \tau_{\mathrm{b}}}\left[x_{k} e^{2 \alpha \tau_{\mathrm{a}}} \cos \beta \tau_{\mathrm{b}}\right. \\
+\frac{1}{\beta}\left\{\alpha x_{k} e^{2 \alpha \tau_{\mathrm{a}}}\right. \\
\left.\left.+\left(\alpha^{2}+\beta^{2}\right) y_{\mathrm{ref}}\right\} \sin \beta \tau_{\mathrm{b}}\right]
\end{array}\right], y_{k}>D} \\
e^{\alpha \tau_{\mathrm{b}}}\left[y_{\mathrm{ref}} \cos \beta \tau_{\mathrm{b}}+\frac{1}{\beta}\left\{x_{k} e^{2 \alpha \tau_{\mathrm{a}}}\right.\right. \\
\left.\left.+\alpha y_{\mathrm{ref}}\right\} \sin \beta \tau_{\mathrm{b}}\right]
\end{array}\right]
\end{align*}
$$

where $\alpha=-\sigma, \beta=\sqrt{\sigma^{2}-1}$, and $\tau_{\mathrm{a}}\left(y_{k}\right)=\left(y_{\text {ref }}-y_{k}\right) / B$, respectively. $\tau_{\mathrm{b}}$ is expressed as follows:

$$
\begin{equation*}
\tau_{\mathrm{b}}\left(y_{k}\right)=T\left[1-\left(\frac{\tau_{\mathrm{a}}}{T} \bmod 1\right)\right] \tag{19}
\end{equation*}
$$



Figure 2: Buck-boost converter.


Figure 3: Behavior of the buck-boost converter.

In addition, the border $D$ is given by

$$
\begin{equation*}
D=y_{\mathrm{ref}}-B T \tag{20}
\end{equation*}
$$

Note that we omit the detailed expressions for the behavior of the waveforms and for the derivation of the Poincare map because they have reported in many previous papers [9,10].

In the following analysis, we fix the parameters $\sigma$ and $T$ as $\sigma=0.2795$ and $T=1.1180$ [11]. Figure 4 shows an example of the 1 parameter bifurcation diagram upon varying the bifurcation parameter $B$ from $B=0.0$ to $B=0.1$. We observe the period-doubling bifurcation and the border-collision bifurcation in the figure. For example, the period- 1 solution bifurcates to the period- 2 solution around $B=0.08$ through the period-doubling bifurcation. On the other hand, it is observed that the border-collision bifurcation occurs and the period-2 solution bifurcates to the period- 4 solution around $B=0.05$; because a part of the period-2 solution collides with the border $D$ and period-4 solution is generated. Specifically, Fig. 5 shows the examples of the waveforms and the corresponding phase plane. This figure corresponds to the parameters (a) and (b) in Fig. 4. It is also clear from Fig. 5 that a part of the period-2 solution bifurcates to the border $D$ and the period- 4 solution is generated. So, we focused on this border-collision bifurcation, and applied the proposed algorithm. Table 1 shows the calculation results. The table says that the border-collision bifurcation appears around $B=0.052$ when $y_{\mathrm{ref}}=0.08$. The bifurcation point $B=0.052$ is correct in the 1 parameter bifurcation diagram (see 4). Therefore, we conclude that the proposed algorithm can calculate the bifurcation point of the border-collision.

## 4. Conclusion

In this paper, we have presented an algorithm for calculating the bifurcation point of the border-collision in the piecewise linear system based on the Poincare map approach. First, we explained the system dynamics. Then, we constructed the algorithm with the exact solution. Finally, we applied the algorithm for the current mode controlled buck-boost converter and calculated the bifurcation point of the border-collision. The proposed algorithm do not need to calculate the Poincaré map and its derivative numerically. We know that the numerical computation processes for deriving the Poincaré map and its derivative is complicated. Therefore, we believe that the proposed algorithm is valuable to calculate the bifurcation point of the bordercollision in the high-dimensional system. Moreover, the mathematical tool MAPLE will help us to derive the exact Poincaré map and its derivative in the high-dimensional systems. In future, we will apply the algorithm to the highdimensional systems.

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Figure 4: 1 parameter bifurcation diagram $\left(y_{\mathrm{ref}}=0.08\right)$.


Figure 5: Waveforms and phase portraits $\left(y_{\mathrm{ref}}=0.08\right)$.

Table 1: Calculation results ( $\sigma=0.2795, T=1.1180$ ).

| $x$ | $y$ | $B$ | $y_{\text {ref }}$ | Remarks |
| :---: | :---: | :---: | :---: | :---: |
| 0.0660 | 0.0246 | 0.0585 | 0.090 | BC |
| 0.0623 | 0.0232 | 0.0552 | 0.085 | BC |
| 0.0587 | 0.0219 | 0.0520 | 0.080 | BC |
| 0.0550 | 0.0205 | 0.0487 | 0.075 | BC |
| 0.0513 | 0.0191 | 0.0455 | 0.070 | BC |

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