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STUDY OF LIGHT BEAM PROPAGATION IN THE ATMOSPHERE BY LASER-RADAR TECHNIQUES USING RAMAN SCATTERING AND ABSORPTION SPECTROSCOPY

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Introduction

Great interest has been concentrated in high power lasers as sources for both optical radar and communication experiments in the atmosphere. There are a number of propagation problems to be thoroughly studied, including molecular and aerosol scattering, molecular absorption, and turbulence. The purpose of this paper is to discuss the performances of the two specific schemes for studying light beam propagation, (1) laser-Raman radar, and (2) resonance absorption method, designed as single ended optical transmissometers.

Laser radar techniques have so far been based on detecting elastically backscattered light which is a combination of Rayleigh scattering by air molecules and Mie scattering by particulate matter such as aerosols and water droplets. Some of the photons in the laser beam are inelastically backscattered at shifted frequencies due to Raman scattering by atmospheric molecules. Because the Raman scattered light from individual constituents is shifted by proper frequency intervals, a means is thereby provided for isolating the scattered light of each of various atmospheric species, and for understanding the nature of almost all molecular constituents.

The laser-Raman radar scheme has been proposed¹ and confirmed experimentally²⁻⁷ to measure simultaneously the composition and the number density of the atmospheric molecules. The resonance absorption scheme has also been discussed previously⁸ to detect small concentration of atmospheric pollutants.

We consider here some potentialities of these two optical schemes in the application to the specific problem of measuring transmission and absorption properties of the laser beam in the atmosphere.

Laser Radar Systems

The laser radar system used in the

experiments basically consists of a laser transmitter and an optical receiver. In the laser-Raman radar, a Q-switched ruby laser at 6943 Å was mainly employed along with a pulsed nitrogen laser at 3371 Å with 50 pps repetition operation. Atmospheric backscatter was collected by a 30-cm diam, 150-cm focal length reflecting telescope. A grating monochromator, as a spectral analyzer, was set to select the Raman-shifted components. The output of the monochromator was detected by a photomultiplier, and then displayed on a 30-MHz oscilloscope. For the resonance absorption scheme, the above mentioned transmitter was replaced by a frequency-tunable dye laser, pumped by a nitrogen laser pulse.

Raman Backscattering Scheme

The Raman backscattered power from atmospheric molecules at a range R is given by the laser-Raman radar equation:

$$P_T = W_0 S T(\omega_b) T(\omega_R) N(R) \sigma_R / R^2 \quad (1)$$

where W_0 is laser output energy, S is system sensitivity, $T(\omega_b)$ and $T(\omega_R)$ are transmittance of the laser light at the frequency of ω_b and of the Raman echo at ω_R , $N(R)$ is concentration of specific molecule, and σ_R is Raman backscattering cross-section of the molecule.

Fig. 1 shows typical example of the oscilloscope traces obtained in the

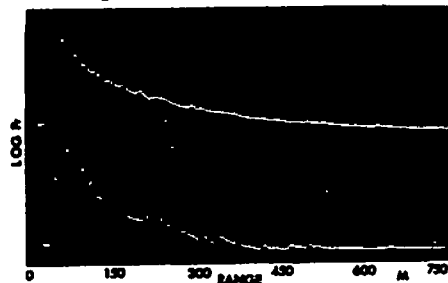


Fig. 1 Laser-Raman echo of O₂ and Mie echo in the clear atmosphere.

