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# Polarization independent microring resonator filter using internal stress and temperature control

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*Abstract* The polarization dependence of resonant wavelength of vertically coupled microring resonator filter was eliminated using the photo-elastic effect induced by the internal stress and the temperature control.

### 1. Introduction

To realize a polarization independent waveguide filter, two methods were proposed. One is a method using half-wave plate inserted in the device, which was adopted for the arrayed waveguide grating filter<sup>[1]</sup>. However, this method is difficult to be applied to ultra compact photonic devices such as microring resonators with several tens micron radius, because the thickness of  $\lambda/2$  plate is ten micron. Another method is the control of aspect ratio of waveguide cross section. However, this method is also difficult to be applied to high index contrast waveguide devices, such as microring resonators, due to their small tolerance.

In this report, we propose and demonstrate the third method using the control of temperature, which is applicable to high index contrast waveguide devices. In addition, we can compensate for the residual polarization dependence of resonant wavelength due to the fabrication error by temperature control. Therefore, polarization independent ultra compact device can be realized easily by this method.

### 2. Temperature dependence control of resonant wavelength by internal stress control

We have reported an athermal waveguide filter<sup>[2]</sup> using polymer as cladding material. Last year, we experimentally realized an athermal vertically coupled microring resonator as shown in Fig. 1 by the control of internal stress as a novel athermal waveguide filter without using polymers<sup>[3]</sup>. This athermalization method is based on the compensation of temperature coefficient of refractive index by their photo-elastic effect induced by the difference of thermal expansion coefficient between the waveguide core and the substrate.

The photo-elastic constant of most silica-based glass materials is negative. In addition, the thermal expansion coefficient of Ta<sub>2</sub>O<sub>5</sub>-SiO<sub>2</sub> compound glass, which was used as the core material in our device<sup>[3]</sup>, ranges from  $0.55 \times 10^{-6}$  [1/K] to  $0.80 \times 10^{-6}$  [1/K] depending on the Ta<sub>2</sub>O<sub>5</sub>-SiO<sub>2</sub> content, which is smaller

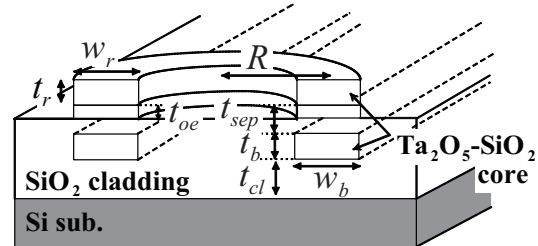


Fig.1 Structure of vertically coupled microring resonator filter.

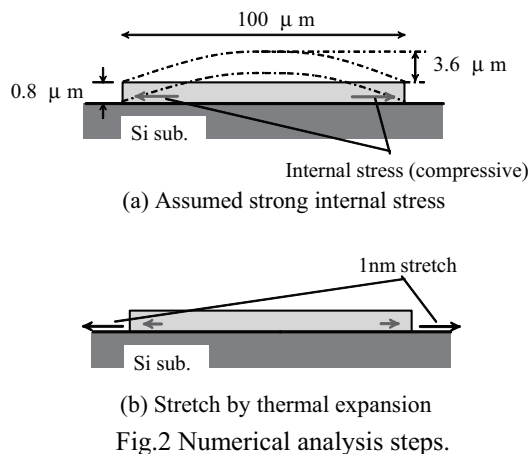
than that of Si substrate ( $2.63 \times 10^{-5}$  [1/K]). Therefore, if the core layer is formed on a Si substrate, the stress of core layer increases with the increase of temperature due to the difference of thermal expansion coefficient between the core and the substrate.

When the internal stress is negligibly small, the photo-elastic effect can be ignored and the temperature coefficient of resonant wavelength  $d\lambda_0/dT$  can be calculated in terms of the temperature coefficient of refractive index of waveguide materials and the thermal expansion coefficient of substrate<sup>[2]</sup>.

On the other hand, when the internal stress is too strong to ignore, the refractive index of core layer suffers the influence of both the thermo-optic effect and the photo-elastic effect. However, the controllability of the temperature dependence by the internal stress control was not made clear. Therefore, we numerically analyzed the photo-elastic effect under a strong internal stress using a finite-element solver (ANSYS by ANSYS, Inc.).

### 3. Numerical investigation of photo-elastic effect

According to our preliminary experiment, a Ta<sub>2</sub>O<sub>5</sub>-SiO<sub>2</sub> straddle mounted beam structure with 100 $\mu$ m length fabricated by RF sputtering deposition was largely deflected by an internal stress, and the deflection was about 3.6 $\mu$ m. The internal stress depends on the sputtering condition. We numerically analyzed the thermo-optic effect of silica-based core layer enhanced by the strong internal stress in the following way.



First, the 100 $\mu\text{m}$  long beam structure was assumed to suffer a strong internal stress which deflected the beam by 3.6 $\mu\text{m}$  (Fig. 2 (a)). When the beam is kept to contact the substrate, a strong stress of the order of  $10^9$  [Pa] of magnitude is induced to the beam at this time. Next, both the bottom edges are assumed to be stretched by 1nm which corresponds to the thermal expansion of substrate by the temperature increase of 10 $^\circ\text{C}$  (Fig. 2 (b)). In this occasion, the internal stress is decreased by the order of  $10^8$  [Pa] of magnitude due to the thermal expansion. The change of the internal stress of a few  $10^8$  [Pa] can result in the change of refractive index by the order of  $10^{-5}$  of magnitude per degree<sup>[4]</sup>, which almost coincides with the experimental result<sup>[3]</sup>. The thermo-optic coefficients of TE and TM modes are calculated to be  $1.269 \times 10^{-5}$  [1/K] and  $1.053 \times 10^{-5}$  [1/K], respectively. It was found from this analysis that the photo-elastic effect can be as large as the thermo-optic effect under a strong internal stress, so as to compensate for the temperature dependence of refractive index.

We also analyzed numerically the photo-elastic effect for the case of no intrinsic internal stress. As a result, when both the bottom edges of the beam were expanded in the same way, the change of internal stress was much smaller by 1/100 than the previous case of strong intrinsic internal stress. Therefore, the change of refractive index is negligibly small without the strong intrinsic internal stress.

#### 4. Elimination of polarization dependence

The temperature coefficient of resonant wavelength of the filter with internal stress has a large polarization dependence<sup>[3]</sup> due to the stress birefringence, compared with conventional microring resonator filters. Therefore, the resonant wavelengths for both polarizations can be equalized at a certain temperature. Fig. 3 shows the temperature dependence of resonant wavelengths of a device with the core width and thickness of 1.4  $\mu\text{m}$  and 0.8  $\mu\text{m}$ , respectively. It is seen that the resonant wavelength for TE and TM polarizations can be

equalized at around 70 $^\circ\text{C}$ . By the precise temperature control, the resonant wavelengths were equalized at 69 $^\circ\text{C}$  as shown Fig. 4, and a polarization independent filter was successfully realized.

When the spectrum peak of TE polarization is athermalized as realized in Ref. [3], the polarization independent filter can be realized by the temperature control without the wavelength shift of TE peak.

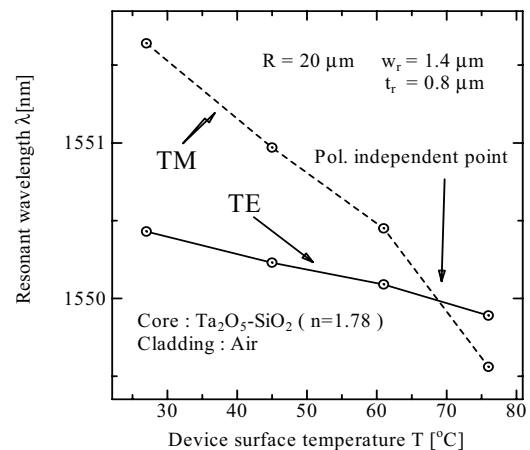


Fig.3 Relation between resonant wavelength and device surface temperature.

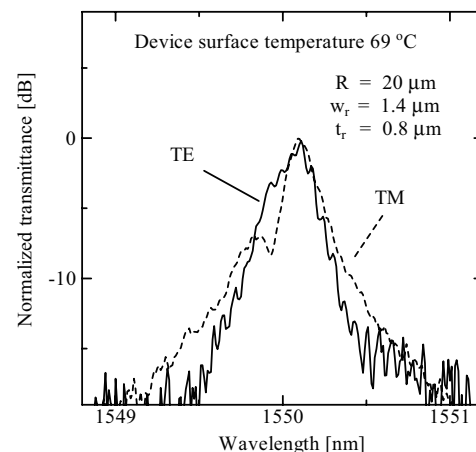


Fig.4 Spectrum response of polarization independent microring resonator filter realized at  $T=69^\circ\text{C}$ .

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