

Basins of Attraction of Steady Operating Conditions in a Two-site Electricity and Heat Supply System

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Abstract—This paper analyzes the basin of attraction of a stable equilibrium point representing a steady operating condition of synchronous generators in a two-site electricity and heat supply system. The analysis is used for considering the effect of heat transfer management on the dynamics of the generators. The basin of attraction becomes small depending on the heat transfer rate, and a change of the setpoints of the combined heat and power plants for regulating the heat transfer rate possibly destabilizes the generators.

1. Introduction

This paper numerically studies a dynamical model of two-site electricity and heat supply system based on our previous studies [1-3]. We examine the basin of attraction of a stable equilibrium point representing a steady operating condition of synchronous generators in the electric sub-system. This is of basic significance for understanding the system's response to an open-loop control of the energy flows in the two-site system. In [3], we proposed a state-feedback (closed-loop) controller that enables regulation of electricity and heat flows based on the nonlinear control technique [4,5]. This controller provides a trajectory of state variable which realizes the desired energy flows. However, it seems not easy to estimate the basin of attraction of the desired trajectory because of the complexity of the closed-loop system, in which the dynamics of the electric and heat sub-systems are coupled each other.

In this paper, we consider an open-loop control in which the set-points of Combined Heat and Power (CHP) plants (including gas turbines and generators) are already determined to realize the desired energy flows. To the open-loop control, the responses of the electric and heat sub-systems can be considered separately. In a viewpoint of dynamical systems theory, the studied model of the electric subsystem appears as a model of double swing dynamics with external forcing [6–8]. In [6–8], the basin structure of the model of swing dynamics was investigated by taking systematic slices of the phase space. In [9], a similar method for analyzing basin structure of dynamical systems is developed by using cell state space and mapping on it. Based on these studies, in [2], the basin portraits of the dynami-

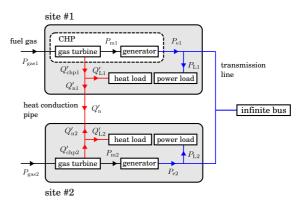


Figure 1: Block diagram of the two-site system. The arrows show the positive directions of energy flows.

cal model of the two-site system was visualized under several fixed values of the set-points of the CHP plants. The visualization was then used for understanding an ideal response of the two-site system to a step-wise change of the set-points of the CHP plants. Here, we discuss this in a more realistic situation, and consider a ramp-wise change of the set-points.

2. Mathematical model

This section introduces a dynamical model of the electricity and heat supply system based on [1, 2]. Figure 1 shows the block diagram of the two-site system in which the positive directions of energy flows are denoted. The notion of site is a unit of energy system that includes a CHP plant, power load, and heat load. The two sites are connected to an infinite bus through a transmission line and are interconnected by a heat conduction pipe.

2.1. Electric sub-system

The electric sub-system in Fig. 1 consists of the two generators, power loads, transmission lines, and infinite bus. The model of electric sub-system is based on the swing equation [10] with δ_i representing the electric angular position of rotor with respect to the infinite bus, and ω_i the deviation of rotor speed relative to the synchronous speed ω_s . The variable δ_i is in the electrical radian, and ω_i is scaled by $\omega_r := \sqrt{\omega_s/2H_i}$, where H_i stands for the per-unit time constant of rotor. The dynamics of generators are represented as follows: for i = 1, 2,

$$\frac{\mathrm{d}\delta_i}{\mathrm{d}t} = \omega_i, \quad \frac{\mathrm{d}\omega_i}{\mathrm{d}t} = P_{\mathrm{m}i} - D_i\omega_i - P_{\mathrm{e}i}(\delta_1, \,\delta_2), \quad (1)$$

where P_{mi} stands for the mechanical input power to the generator, and D_i for the damping coefficient. The function P_{ei} stands for the electric output power and is given by

$$P_{ei} = \sum_{j \in \{1,2,\infty\}} E_i E_j \{G_{ij} \cos(\delta_i - \delta_j) + B_{ij} \sin(\delta_i - \delta_j)\}, \quad (2)$$

with the symbol ∞ representing the infinite bus, and $\delta_{\infty} = 0$. The parameter E_i corresponds to the voltage behind synchronous reactance, and $G_{ij} + iB_{ij}$ are the transfer admittances.

2.2. Heat sub-system

The heat sub-system in Fig. 1 consists of the conduction pipe and loads. Here, we do not consider the transient dynamics and losses of heat transfer through the heat conduction pipe. This is relevant for considering the open-loop control of the two-site system. By using the following model, the set-points, i.e. the fuel inputs to the CHP plants, are determined to realize a desired heat transfer rate Q'_n . In Fig. 1, the conservation of energy at each site induces the following equality:

$$Q'_{\rm chpi} = Q'_{\rm ni} + Q'_{\rm Li}.$$
 (3)

Further, the heat output rates Q'_{n1} and Q'_{n2} satisfy

$$Q'_{n1} = -Q'_{n2} := Q'_n, \tag{4}$$

where Q'_n represents the heat transfer rate from site #1 to site #2.

2.3. Gas turbine

The gas turbine at site #*i* converts the gas input rate $P_{\text{gas}i}$ to both the mechanical power $P_{\text{m}i}$ and the heat rate $Q'_{\text{chp}i}$. Because its time response is sufficiently fast compared with the electromechanical dynamics of the generators [11], the dynamics of the gas turbine are not considered in this paper. Then, the instantaneous conversion of energy at each gas turbine is represented by

$$\begin{bmatrix} P_{mi} \\ Q'_{chpi} \end{bmatrix} = \begin{bmatrix} \eta_{ei} \\ \eta_{hi} \end{bmatrix} P_{gasi}.$$
 (5)

Throughout this paper, the parameters η_{ei} and η_{hi} are constant and satisfy $\eta_{ei} + \eta_{hi} < 1$. The constant η_{ei} represents the thermal efficiency of the gas turbine at site #i, and η_{hi} the ratio of heat output rate to gas input rate.

2.4. Derived model

Consequently, the dynamics of the two-site electricity and heat supply system are represented by the following nonlinear dynamical model:

$$\frac{\mathrm{d}\delta_1}{\mathrm{d}t} = \omega_1,\tag{6a}$$

$$\frac{d\omega_1}{dt} = \frac{\eta_{e1}}{\eta_{h1}} (Q'_n + Q'_{L1}) - D_1 \omega_1 - P_{e1}(\delta_1, \delta_2), \qquad (6b)$$

$$\frac{\mathrm{d}\delta_2}{\mathrm{d}t} = \omega_2,\tag{6c}$$

$$\frac{d\omega_2}{dt} = \frac{\eta_{e2}}{\eta_{h2}} (-Q'_n + Q'_{L2}) - D_2\omega_2 - P_{e2}(\delta_1, \,\delta_2).$$
(6d)

The dynamical model (6) contains the parameters Q'_n and Q'_{Li} of the heat sub-system. In the rest of this paper, with this model, the effect of the heat sub-system on dynamics of the electric sub-system will be studied.

3. Steady operating conditions

This section analyzes equilibrium points of the dynamical model (6) in order to investigate the steady operating conditions of the generators. Since the dynamical model (6) has the same formulation as the classical swing equations, the analysis method used in [12] is applied for investigating how the steady state characteristics depend on Q'_n . From the condition $d\delta_i/dt = 0$ at equilibrium points, we have

$$\omega_i^* = 0, \tag{7}$$

where ω_i^* represents the value of ω_i at equilibrium points. From the condition $d\omega_i/dt = 0$, the values of phase angles δ_1^* and δ_2^* satisfy the following equations:

$$\alpha_{1} = \sin \delta_{1}^{*} + \kappa_{1} \sin(\delta_{1}^{*} - \delta_{2}^{*}) + \lambda_{1} \cos \delta_{1}^{*} + \mu_{1} \cos(\delta_{1}^{*} - \delta_{2}^{*}),$$

$$\alpha_{2} = \sin \delta_{2}^{*} + \kappa_{2} \sin(\delta_{2}^{*} - \delta_{1}^{*}) + \lambda_{2} \cos \delta_{2}^{*} + \mu_{2} \cos(\delta_{2}^{*} - \delta_{1}^{*}),$$
(8)

where α_1 and α_2 are defined by

$$\alpha_1 := \frac{\eta_{e1}(Q'_{L1} + Q'_n) - \eta_{h1}E_1^2G_{11}}{\eta_{h1}E_1E_{\infty}B_{1\infty}},$$
(9a)

$$\alpha_2 := \frac{\eta_{e2}(Q'_{L2} - Q'_n) - \eta_{h2}E_2^2 G_{22}}{\eta_{h2}E_2 E_\infty B_{2\infty}},$$
(9b)

and κ_i , λ_i , and μ_i are given by

$$\kappa_i = \frac{E_1 E_2 B_{12}}{E_i E_\infty B_{i\infty}}, \quad \lambda_i = \frac{G_{i\infty}}{B_{i\infty}}, \quad \mu_i = \frac{E_1 E_2 G_{12}}{E_i E_\infty B_{i\infty}}.$$
 (10)

By solving the equation (8), the values of δ_1^* and δ_2^* are numerically determined. Fig. 2 shows the result on existence and number of equilibrium points. The values of parameters are shown in Tab. 1. In the region R_n (n = 2, 4, 6), there are n distinct equilibrium points. In the three regions, one of the equilibrium points is asymptotically stable, and the others are unstable. The stable equilibrium point represents a synchronized motion of the two generators in which they operate with the same frequency as the infinite bus.

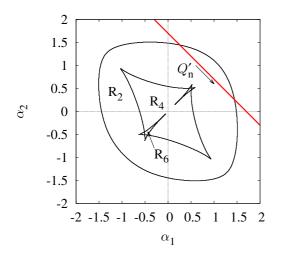


Figure 2: Numerical result on existence and number of equilibrium points. In the region R_2 (or R_4 , R_6), there are two (or four, six) distinct equilibrium points of the dynamical model (6).

Table 1: Values of parameters for numerical analysis

| Rated power | $P_{\rm b}$ | 1.0 MW |
|---|----------------------|-----------------------------|
| Synchronous speed | $\omega_{ m b}$ | $2\pi \cdot 60 \mathrm{Hz}$ |
| Inertial constant | H_i | 10 s |
| Damping coefficient | D_i | 0.005 |
| Voltage | E_i | 1.0 |
| Transfer susceptance $(\#i, \infty)$ | $B_{i\infty}$ | 1.0 |
| Transfer conductance $(\#i, \infty)$ | $G_{i\infty}$ | -0.1 |
| Transfer susceptance (#1, #2) | B_{12} | 0.5 |
| Transfer conductance (#1, #2) | G_{12} | 0.05 |
| Transfer conductance (# <i>i</i> , # <i>i</i>) | G_{ii} | 0.05 |
| Heat load | $Q_{\mathrm{L}i}$ | 0.9 |
| Coefficient of electricity output | $\eta_{{ m e}i}$ | 0.40 |
| Coefficient of heat output | $\eta_{\mathrm{h}i}$ | 0.40 |

Here, we consider the stability of the equilibrium points due to the quasi-static changes of the parameters of the heat sub-system. As Q'_n changes, the steady operating point moves in the (α_1, α_2) -plane along the straight line given by

$$e_1\alpha_1 + e_2\alpha_2 = (Q'_{L1} + Q'_{L2}) - e_3 \tag{11}$$

where the coefficients e_1 to e_3 are determined by the parameters of the electric sub-system and are given by

$$e_{1} := \frac{\eta_{h1}}{\eta_{e1}} E_{1} E_{\infty} B_{1\infty}, \quad e_{2} := \frac{\eta_{h2}}{\eta_{e2}} E_{2} E_{\infty} B_{2\infty},$$

$$e_{3} := \frac{\eta_{h1}}{\eta_{e1}} E_{1}^{2} G_{11} + \frac{\eta_{h2}}{\eta_{e2}} E_{2}^{2} G_{22}.$$
(12)

The equation (11) is obtained by eliminating Q'_n from (9). Since the line (11) is parameterized by $Q'_{L1} + Q'_{L2}$, a steady operating condition is determined by the values of Q'_n and $Q'_{sum} := Q'_{L1} + Q'_{L2}$. In Fig. 2, the *red* line shows (11) with $Q'_{sum} = 1.8$. The synchronized operation of the generators is achieved when the operating condition determined by Q'_n and Q'_{sum} is kept within R₂.

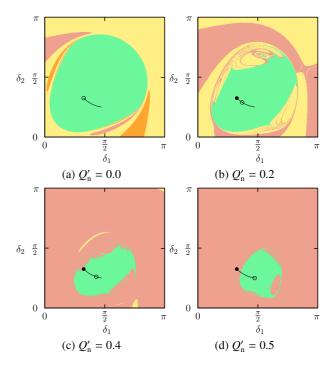


Figure 3: Visualization of basins of attraction. The *solid* line shows the equilibrium points under various Q'_n between 0 and 0.5. The *circle* (\circ) shows the equilibrium point for each Q'_n , and the dot (\bullet) for $Q'_n = 0$.

4. Basins of attraction

This section analyzes the basins of attraction of the stable equilibrium points under several fixed values of Q'_n . Based on the analysis, we consider the effect of heat transfer management on the dynamics of the electric subsystem. A possible open-loop control is then discussed in terms of a transient instability. Following [6–8], the basin of attraction is visualized by taking a two-dimensional slice of $\{(\delta_1, \delta_2, \omega_1, \omega_2) \in \mathbb{X} \mid \omega_1 = 0, \omega_2 = 0\}$ in the entire phase space $\mathbb{X} := \mathbb{T}^2 \times \mathbb{R}^2$, where \mathbb{T} stands for the torus, and \mathbb{R} for the set of real number. For the slice, initial conditions on a grid of 401×401 points were numerically integrated. Each point is colored according to the attractor reached from the corresponding initial condition.

Fig. 3 shows the visualization of the basins of attraction under several values of Q'_n . Under the current setting of the parameters, the system (6) has four attractors. One attractor is the stable equilibrium point representing the steady operating condition: this is shown by *circle* (\circ) in the figure, and its basin is colored *green*. A second attractor is a periodic orbit, in which the generator #1 operates at a desynchronized manner with the infinite bus: its basin is colored *red*. A third one is another periodic orbit, in which the generator #2 is desynchronized; its basin is colored *orange*. In the forth attractor, both generators are desynchronized; its basin is *yellow*. Fig. 3 indicates that the heat transfer management affects the responses of the electric sub-system, and the basins of attraction of the stable

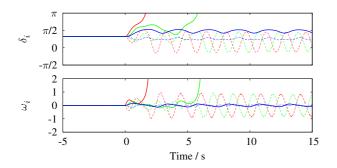


Figure 4: System's responses to ramp-wise changes of the heat transfer rate under $T_d = 0$, 0.5, and 1.0 s. The *solid* line shows the response of the generator #1, and the *broken* line the generator #2.

equilibrium points become small on the two-dimensional slices as Q'_n increases. In [2], a similar result is obtained for $D_i = 0.21$ and $\lambda_i = \mu_i = 0$.

This analysis suggests a possibility of instability due to a change of Q'_n . In Fig. 3, the *solid* line shows the stable equilibrium points under various Q'_n between 0 and 0.5, and the dot (•) the equilibrium point under $Q'_n = 0$. The basins of attraction directly illustrates the following two ideal operations of the two-site system. Since there exists an equilibrium point for each Q'_n , quasi-static change of the set-points of the CHP plants enables the change of operating conditions of the generators along the lines in Fig. 3. However, a step-wise change of Q'_n from 0 to 5.0 desynchronizes the generator #1 because the dot (•) exists outside the domain of attraction of the stable equilibrium point in Fig. 3d.

As a realistic situation, an open-loop control of the heat transfer rate Q'_n can be considered as in between the above two ideal situations. In this paper, based on [13], we consider a ramp-wise change of the set-points of CHP plants from $Q'_{\rm n} = 0$ to 0.5. The duration $T_{\rm d}$ of the change of the set-points is an important parameter: $T_d = 0$ corresponds to the step-wise change, and $T_{\rm d} = \infty$ the quasistatic change. In an engineering viewpoint, the range of T_d where the instability does not occur is of significant importance. Fig. 4 shows the system's responses for $T_d = 0, 0.5$, and 1.0 s. The *red* line shows the case of $T_d = 0$ (stepwise change), and the generator #1 is desynchronized as mentioned above. In the case of $T_d = 1.0 \text{ s}$ (*blue* line), the variables δ_i and ω_i converged to the values of the equilibrium point. In the case of $T_d = 0.5$ s, it is observed that the generator #1 is desynchronized. The analysis of the relationship between T_d and the basins of attraction is future work and discussed at the end of this paper.

5. Conclusions and discussion

In this paper, we analyzed the basins of attraction of equilibrium points representing steady operating conditions of synchronous generators in a two-site electricity and heat supply system. The slices of the basins were visualized under various fixed values of heat transfer rate Q'_n .

The analysis indicated that the heat transfer management affected the responses of the electric sub-system, and the basin of attraction of the stable equilibrium point became small depending on Q'_n . Furthermore, a possibility of instability was discussed for step-wise and ramp-wise changes of the set-points of the CHP plants. It was observed that the instability possibly occurred under a ramp-wise change with a small values of the duration T_d .

Finally, for the future work, we discuss the possibility of analyzing T_d via the basins of attraction. After the time $t = T_d$, from the uniqueness of the solution of (6), the resultant behavior is determined by the basins in Fig. 3 if the state trajectory passes the slice determined by $\omega_1 = \omega_2 = 0$. However, in general, this is not the case because the two dimensional slice is not transversal in the full four-dimensional phase space. Nevertheless, it is observed in the case of $T_d = 0.5$ s (green line) in Fig. 4 that the state passes a slice of $\{(\delta_1, \delta_2, \omega_1, \omega_2) \in \mathbb{X} \mid \omega_1 = \epsilon, \omega_2 = 0\}$ for a small ϵ . Thus, if the basin structure does not vary drastically depending on the values of ω_i , the visualization of the basins in Fig. 3 may be used for the analysis.

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